# Updates on junctions in PYTHIA and describing strangeness 

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## Hadronisation and the Lund String Model

Hadronisation in PYTHIA:
$>$ Maps partons to hadrons using the Lund String Model
> Represent the colour-confinement field between colourconnected partons (i.e. form overall colour singlet state) as strings
$>$ Partons move apart and "break" the string, creating new light quark-antiquark pairs (or diquark-antidiquark pairs)


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Starting point is Leading Colour limit $N_{C} \rightarrow \infty$
$>$ Each colour is unique $\rightarrow$ only one way to make colour singlets

In $e^{+} e^{-}$collisions (LEP):
$>$ Corrections suppressed by $1 / N_{C}^{2} \sim 10 \%$
$>$ Not much overlap in phase space


e.g. a dipole string configuration which make
use of the colour-anticolour singlet state

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Starting point is Leading Colour limit $N_{C} \rightarrow \infty$
$>$ Each colour is unique $\rightarrow$ only one way to make colour singlets

But high-energy pp collisions involve very many coloured partons with significant phase space overlaps

e.g. a dipole string configuration which make use of the colour-anticolour singlet state

## QCD Colour Reconnections

Stochastically restores colour-space ambiguities according to SU(3) algebra
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## Dipole-type reconnection

What about the
-green-blue colour singlet state?


Junctions!

## Junctions



## Junctions



Mechanism for baryon production
$>\sim 40 \%$ of baryons are from junctions in PYTHIA

## Asymmetries

$>$ Equal amount of junctions and anti junctions are formed Junctions typically form between jets $\rightarrow$ as jets are likely to have large opening angles due to available phase space, junction sits at low $p_{\perp}$


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Junctions typically form between jets $\rightarrow$ as jets are likely to have large opening angles due to available phase space, junction sits at low $p_{\perp}$

Heavy flavour ba $>\sim 70 \%$ of hea

## Current implementation

$>$ Runs into cases with no solution (particularly for heavy quarks)
$>$ Relies on convergence procedure that fails $\sim 10 \%$ of the time


## Junction Rest Frame

## What is the junction rest frame?

If the momenta of the junction legs are at $120^{\circ}$ angles
$\rightarrow$ the pull in each direction on the junction is equal
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Does a boost to the mercedes frame always exist?

Consider the following:
In the rest frame of one of the partons, and the angle between the other two partons is greater than $120^{\circ}$
*no special consideration for these cases in current implementation


## Pearl-on-a-string

The junction gets "stuck" to the soft quark, which we call a pearl-on-a-string
> More likely to occur for junctions with heavy flavour endpoints

Example of pearl-on-a-string viewed in the Ariadne frame of the green quark


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For a junction to make a heavy baryon, the junction leg with the heavy quark can't fragment (i.e. a "soft" junction leg) = pearl-on-a-string!


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$>$ Average over the pearl motion

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How do we fragment pearl-on-a-string cases?
$>$ Average over the pearl motion
$>$ Fragment like a $q-g-\bar{q}$ string typically only a good approximation for light quarks


## Updates to averaging



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Use an "average" JRF
Current procedure assumes the average is the mercedes frame
$>$ Uses energy weighted sum of momenta on each junction leg
$>$ Relies on convergence procedure that fails $\sim 10 \%$ of cases

New treatment:
$>$ Considers pull on junction over time and average over junction motion

- Includes pearl-on-a-string
$>$ Allow endpoint oscillations
$>$ No reliance on convergence
$>$ Early time JRF defined by the first parton on each leg
$>$ Use smallest leg momentum as a measure of effective time for the JRF
$>$ When softest parton has lost its momentum, the next parton dominates the pull


## Junctions



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## Strangeness Enhancement



Clear observations of strangeness enhancement with respect to charged multiplicity [e.g. ALICE Nature Pays. 13,535 (2017)]

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## Close-packing



Dense string environments
$\rightarrow$ Casimir scaling of effective string tension
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Diquark formation via successive colour fluctuations


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## Thank you for listening!

## Backup Slides

## Junction Rest Frame

## What is the junction rest frame?

The standard JRF for a 3-parton configuration is defined as when the angle between each legs 3 -momentum is $120^{\circ}$

## Finding the $120^{\circ}$ JRF:

$>$ Fix the angle between the 3-momenta to $120^{\circ}$ and calculate the invariants:

$$
a_{i j}=p_{i}^{\prime} p_{j}^{\prime}=E_{i}^{\prime} E_{j}^{\prime}-\left|\vec{p}_{i}^{\prime}\right|\left|\vec{p}_{j}^{\prime}\right| \cos \frac{2 \pi}{3}=E_{i}^{\prime} E_{j}^{\prime}+\frac{1}{2}\left|\vec{p}_{i}^{\prime}\right|\left|\vec{p}_{j}^{\prime}\right| \quad \text { Where dashed values are energy and momentum in the } 120^{\circ} \mathrm{JRF}
$$

$>$ Therefore we can make equation

$$
f_{i j}=f\left(\left|{\overrightarrow{p^{\prime}}}_{i}^{\prime}\right|,\left|\vec{p}_{j}^{\prime}\right| ; m_{i}, m_{j}, a_{i j}\right)=\sqrt{\left|\vec{p}_{i}^{\prime}\right|^{2}+m_{i}^{2}} \sqrt{\left|\vec{p}_{j}^{\prime}\right|^{2}+m_{j}^{2}}+\frac{1}{2}\left|\vec{p}_{i}^{\prime}\right|\left|\vec{p}_{j}^{\prime}\right|-a_{i j}
$$

with solutions when $f_{i j}=0$
$>$ Set $f_{13}=f_{12}=0$, and solving for $\left|\vec{p}_{j}\right|$ in terms of $\left|\vec{p}_{1}\right|$ we get

$$
\left|\vec{p}_{j}^{\prime}\right|\left(\vec{p}_{1}^{\prime}\right)=\frac{2 E_{1}^{\prime} \sqrt{4 a_{1 j}^{2}-m_{j}^{2}\left(4 E_{1}^{\prime 2}-\left|\vec{p}_{1}^{\prime}\right|^{2}\right)}-2\left|\vec{p}_{1}^{\prime}\right| a_{i j}}{4 E_{1}^{\prime 2}-\left|\vec{p}_{1}^{\prime}\right|^{2}}
$$

Then what if the JRF is no solution to $f_{23}=0$ ?
$>\operatorname{Set} f_{23}=0$, sub in the above equations for $\left|\vec{p}_{2}^{\prime}\right|$ and $\left|\vec{p}_{3}^{\prime}\right|$, and solve for $\left|\vec{p}_{1}^{\prime}\right|$

## Junction Rest Frame

The junction gets "stuck" to the soft quark, which we call a pearl-on-a-string
Consider a basic case:
$>$ Two massless legs and one massive soft leg in the Ariadne frame with respect to the massive parton

$$
\begin{aligned}
p(t) & =p_{0}-2 \kappa x(t)=\frac{m v(t)}{\sqrt{1-v(t)^{2}}} \\
\frac{d x}{d t} & =\frac{1}{\sqrt{1+\frac{m^{2}}{\left(p_{0}-2 \kappa x\right)^{2}}}} \quad x_{\max }=p_{0} / 2 \kappa
\end{aligned}
$$

The differential equation is
non-trivial and not straightforward to compute!!!


## Junction Fragmentation

How do we fragment these junction systems? How do we get the junction baryon?

Use similar fragmentation method as with dipole strings, fragmenting off on-shell hadrons from each junction-leg string end. Treat each junction leg as half a dipole string.

Standard Procedure:
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$>$ Combine partons from last break of two lowest energy strings into a diquark, $q_{3} q_{5}$
$>$ Fragment the last junction leg as dipole with endpoints


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Modelling particularly important for heavy flavour baryons as they are more sensitive to junction motion


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The differential equation is
non-trivial and not straightforward to compute!!!


## Implementation

Know JRF given 3-parton configuration, however what about junctions with gluons? Do not want to map the junction motion and fragmentation in a space-time picture, so instead need some "average JRF" to describe the junction motion.

Current procedure finds average "pull" on junction of each leg and looks for $120^{\circ}$ frame given the average pulls. Problems in current procedure in PYTHIA: $>$ Convergence failure of iterative procedure for about 10\% of junction systems $>$ Only considers $120^{\circ}$ JRF
$>$ No special handling if there is no $120^{\circ}$ frame
$>$ Weightings used in averaging procedure not most physically logical


Look at JRFs at different time steps and average over junction velocities.

## Implementation

## New iterative procedure:

1. Find JRF using the first parton on each junction leg, store the associated velocity, and boost to this frame.
A. If $120^{\circ}$ frame does not exist, use rest frame of soft quark as an approximation of the pearl-on-a-string treatment
(should only occur for massive endpoints)


## Implementation

## New iterative procedure:

1. Find JRF using the first parton on each junction leg, store the associated velocity, and boost to this frame.
2. Time associated with JRF: $p_{\text {small }}=$ smallest absolute 3-momentum
A. If the smallest 3 -momentum is zero, let $p_{\text {small }}$ be the next lowest 3momentum


## Implementation

## New iterative procedure:

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A. If the smallest 3 -momentum is zero, let $p_{\text {small }}$ be the next lowest 3 momentum
3. Pull vectors: Store 4-momenta scaled down (conserving mass) to have 3-momentum magnitude of $p_{\text {small }}$.
A. If at rest, store the rest frame momentum.


## Implementation

## New iterative procedure:

1. Find JRF using the first parton on each junction leg, store the associater velocity, and boost to this frame.
2. Time associated with JRF: $p_{\text {small }}=$ smallest absolute 3-momentum
3. Pull vectors: Store 4-momenta scaled down (conserving mass) to have 3 -momentum magnitude of $p_{\text {small }}$.
4. Update momenta:
A. For small leg
i. Step to next parton on leg if possible.


## Implementation

## New iterative procedure:

1. Find JRF using the first parton on each junction leg, store the associated velocity, and boost to this frame.
2. Time associated with JRF: $p_{\text {small }}=$ smallest absolute 3-momentum
3. Pull vectors: Store 4-momenta scaled down (conserving mass) to have 3 -momentum magnitude of $p_{\text {small }}$.

## 4. Update momenta

5. Find JRF with new 3-parton configuration and iterate: Repeat steps 2 4 till either:

A. the sum of all $p_{\text {small }}$ exceeds 10 GeV
B. two endpoints are reached
C. parton associated with $p_{\text {small }}$ is a massless endpoint.

## Average JRF

## Averaging procedure:

Concerned with the junction motion in the time-frame of the hadronisation process
$\rightarrow$ introduce normalisation parameter $p_{\text {norm }}=2 \mathrm{GeV}$ by default

Expect early time pulls to more heavily influence junction motion
$\rightarrow$ use exponential weighting to model time dependence

$$
v_{j u n}=\frac{\sum_{i=1}^{i_{\max }} v_{i}\left(e^{-p_{i-1} / p_{\text {norm }}^{\prime}}-e^{-p_{i} / p_{\text {norm }}^{\prime}}\right)}{1-e^{-p_{i_{\text {max }}} / p_{\text {norm }}^{\prime}}}
$$

The same averaging procedure is used to calculate the average pull on the junction by each leg
$\rightarrow$ used to construct fictitious endpoints for fragmentation

Mathematical subtleties:
$>$ Each $p_{\text {small }}$ is measured in the successive JRFs, therefore transform by $\gamma$-factor to lab frame
$>p_{\text {norm }}$ is recalculated to consider $\gamma$-factors

$$
p_{\text {norm }}^{\prime}=\sum_{i=1}^{N} \gamma_{i} p_{\text {small }_{i}}+\gamma_{N+1}\left(p_{\text {norm }}-\sum_{i=1}^{N} p_{\text {small }_{i}}\right)
$$

## Pearl-on-a-string

How do we implement pearl-on-a-string model? What is the Ariadne frame if we have gluons the junction legs? Instead, we model the soft quark as a gluon with momentum determined by the average JRF.

Fragment $q-q_{\text {pearl }}-q$ string as a $q-g-\bar{q}$ string using existing fragmentation mechanism in PYTHIA

$>$ Fragment the $q-g_{\text {pearl }}-\bar{q}$ string system from the $\bar{q}$ end, reversing the hadron IDs
$>$ Pick up quark and energy from $p_{\text {pearl }}$ for "free" when stepping over junction


[^0]:    blue $q \bar{q}$ fluctuation breaks nearby blue string, preventing diquark formation

