Updates on junctions in PYTHIA and describing strangeness

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Hadronisation and the Lund String Model

Hadronisation in PYTHIA:

- Maps partons to hadrons using the Lund String Model
- Represent the colour-confinement field between colourconnected partons (i.e. form overall colour singlet state) as strings
- \succ Partons move apart and "break" the string, creating new **light** quark-antiquark pairs (or diquark-antidiquark pairs)







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- Starting point is Leading Colour limit $N_C \rightarrow \infty$
 - \succ Each colour is unique \rightarrow only one way to make colour singlets

In e^+e^- collisions (LEP):

- > Corrections suppressed by $1/N_C^2 \sim 10\%$
- > Not much overlap in phase space



e.g. a dipole string configuration which make use of the **colour-anticolour** singlet state





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But high-energy pp collisions involve very many coloured partons with significant phase space overlaps



e.g. a dipole string configuration which make use of the **colour-anticolour** singlet state

QCD Colour Reconnection (CR) model





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Stochastically restores colour-space ambiguities according to SU(3) algebra

> Allows for reconnections to minimise string lengths











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Heavy flavour baryons





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If the momenta of the junction legs are at 120° angles \rightarrow the pull in each direction on the junction is equal \rightarrow junction is at rest

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Does a boost to the mercedes frame always exist?

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What is the junction rest frame?

If the momenta of the junction legs are at 120° angles \rightarrow the pull in each direction on the junction is equal \rightarrow junction is at rest

*no special consideration for these cases in current implementation

The **junction gets "stuck"** to the soft quark, which we call a **pearl-on-a-string**

More likely to occur for junctions with heavy flavour endpoints

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For a junction to make a **heavy baryon**, the junction leg with the heavy quark can't fragment (*i.e.* a "soft" junction leg) = pearl-on-a-string!

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How do we fragment pearl-on-a-string cases?

> Average over the pearl motion

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endpoints

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- Current procedure assumes the average is the mercedes frame > Uses energy weighted sum of momenta on each junction leg > Relies on convergence procedure that fails ~10% of cases

Updates to averaging

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New treatment:

- > Considers pull on junction over time and average over junction motion
- > Includes pearl-on-a-string
- > Allow endpoint oscillations
- > No reliance on convergence

 \succ Early time JRF defined by the first parton on each leg > Use smallest leg momentum as a measure of effective time for the JRF \succ When softest parton has lost its momentum, the next parton dominates the pull

Updates to averaging

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Clear observations of strangeness enhancement with respect to charged multiplicity [e.g. ALICE Nature Pays. 13, 535 (2017)]

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Dense string environments

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String tension could be different from the vacuum case compared to near a junction

Diquark formation via successive colour fluctuations

Diquark Suppression

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Diquark formation via successive colour fluctuations

blue $q\bar{q}$ fluctuation on the string

Diquark Suppression

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Diquark formation via **successive colour fluctuations**

What if there's a blue string nearby?

Diquark Suppression

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What if there's a blue string nearby?

blue $q\bar{q}$ fluctuation breaks nearby blue string, preventing diquark formation

Diquark Suppression

Diquark formation via **successive colour fluctuations**

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Thank you for listening!

Backup Slides

What is the **junction rest frame**?

The standard JRF for a 3-parton configuration is defined as when the angle between each legs 3-momentum is 120°

Finding the 120° JRF:

> Fix the angle between the 3-momenta to 120° and calculate the invariants: $a_{ij} = p'_i p'_j = E'_i E'_j - |\overrightarrow{p'_i}| |\overrightarrow{p'_j}| \cos \frac{2\pi}{3} = E'_i E'_j + \frac{1}{2} |\overrightarrow{p'_i}| |\overrightarrow{p'_j}| \quad \text{Where dashed values are energy and momentum in the 120° JRF}$ > Therefore we can make equation $f_{ij} = f(|\vec{p'}_i|, |\vec{p'}_j|; m_i, m_j, a_{ij}) = \sqrt{|\vec{p'}_i|^2 + m_i^2} \sqrt{|\vec{p'}_i|^2} + m_i^2 \sqrt{|\vec{p'}_i|^2$ with solutions when $f_{ii} = 0$ > Set $f_{13} = f_{12} = 0$, and solving for $|\overrightarrow{p}_i|$ in terms of $|\overrightarrow{p}_1|$ we get Then what if the JRF is $|\overrightarrow{p'_{j}}|(\overrightarrow{p'_{1}}) = \frac{2E'_{1}\sqrt{4a^{2}_{1j} - m^{2}_{j}(4E'^{2}_{1} - |\overrightarrow{p'_{1}}|^{2})} - 2|\overrightarrow{p'_{1}}|a_{ij}}{4E'^{2}_{1} - |\overrightarrow{p'_{1}}|^{2}}$ no solution to $f_{23} = 0$?

> Set $f_{23} = 0$, sub in the above equations for $|\vec{p}_2|$ and $|\vec{p}_3|$, and solve for $|\vec{p}_1|$

$$\overline{|\vec{p}_{j}'|^{2} + m_{j}^{2}} + \frac{1}{2} |\vec{p}_{i}'| |\vec{p}_{j}'| - a_{ij}$$

The junction gets "stuck" to the soft quark, which we call a pearl-on-a-string

Consider a basic case:

Two massless legs and one massive soft leg in the Ariadne frame with respect to the massive parton

$$p(t) = p_0 - 2\kappa x(t) = \frac{mv(t)}{\sqrt{1 - v(t)^2}}$$
$$\frac{dx}{dt} = \frac{1}{\sqrt{1 + \frac{m^2}{(p_0 - 2\kappa x)^2}}} \quad x_{max} = p_0/2\kappa$$

The differential equation is non-trivial and not straightforward to compute!!!

How do we fragment these junction systems? How do we get the junction baryon?

Use similar fragmentation method as with dipole strings, fragmenting off on-shell hadrons from each junction-leg string end. Treat each junction leg as half a dipole string.

Standard Procedure:

> Go to junction rest frame (JRF)

 q_{02}

 q_{01}

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- > Fragment the last junction leg as dipole with endpoints

 $q_3 q_5 - q_{03}$

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Modelling particularly important for heavy flavour baryons as they are more sensitive to junction motion

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The differential equation is non-trivial and not straightforward to compute!!!

Know JRF given 3-parton configuration, however what about junctions with gluons? Do not want to map the junction motion and fragmentation in a space-time picture, so instead need some "average JRF" to describe the junction motion.

Current procedure finds average "pull" on junction of each leg and looks for 120° frame given the average pulls. **Problems in current procedure in PYTHIA:** > Convergence failure of iterative procedure for about 10% of junction systems > Only considers 120° JRF > No special handling if there is no 120° frame

 \succ Weightings used in averaging procedure not most physically logical

Need more rigorous handling in order to be able to draw solid physics conclusions from the results

Look at JRFs at different time steps and average over junction velocities.

 q_{32}

New iterative procedure:

- 1. Find JRF using the first parton on each junction leg, store the associated velocity, and boost to this frame.
 - A. If 120° frame does not exist, use rest frame of soft quark as an approximation of the pearl-on-a-string treatment

(should only occur for massive endpoints)

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New iterative procedure:

- 1. Find JRF using the first parton on each junction leg, store the associated velocity, and boost to this frame.
- 2. Time associated with JRF: p_{small} = smallest absolute 3-momentum
 - A. If the smallest 3-momentum is zero, let p_{small} be the next lowest 3momentum

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- 3. **Pull vectors:** Store 4-momenta scaled down (conserving mass) to have 3-momentum magnitude of p_{small} .
 - A. If at rest, store the rest frame momentum.

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- 3. **Pull vectors:** Store 4-momenta scaled down (conserving mass) to have 3-momentum magnitude of p_{small} .
- 4. Update momenta:
 - A. For small leg
 - Step to next parton on leg if possible.
 - If massive endpoint, reduce the endpoint to at rest. II.
 - iii. If massless endpoint, make this final iteration.
 - B. Reduce the momentum of the other partons by p_{small} .

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- 3. **Pull vectors:** Store 4-momenta scaled down (conserving mass) to have 3-momentum magnitude of p_{small} .
- 4. Update momenta
- 5. Find **JRF with new 3-parton configuration** and iterate: Repeat steps 2 -4 till either:
 - A. the sum of all p_{small} exceeds 10 GeV
 - B. two endpoints are reached
 - C. parton associated with p_{small} is a massless endpoint.

g₃₁

Averaging procedure:

Concerned with the junction motion in the time-frame of the hadronisation process \rightarrow introduce **normalisation parameter** $p_{norm} = 2$ GeV by default

 \rightarrow use **exponential weighting** to model time dependence

$$v_{jun} = \frac{\sum_{i=1}^{i_{max}} v_i (e^{-p_{i-1}/p'_{norm}} - e^{-p_i/p'_{norm}}}{1 - e^{-p_{i_{max}}/p'_{norm}}}$$

The same averaging procedure is used to calculate the average pull on the junction by each leg → used to construct **fictitious endpoints** for fragmentation

Average JRF

- Expect early time pulls to more heavily influence junction motion

 - norm

Mathematical subtleties:

> Each p_{small} is measured in the successive JRFs, therefore transform by γ -factor to lab frame $> p_{norm}$ is recalculated to consider γ -factors

$$p'_{norm} = \sum_{i=1}^{N} \gamma_i p_{small_i} + \gamma_{N+1} (p_{norm} - \sum_{i=1}^{N} p_{small_i})$$

How do we implement pearl-on-a-string model? What is the **Ariadne frame** if we have gluons the junction legs? Instead, we model the **soft quark as a gluon** with momentum determined by the average JRF.

Fragment $q - q_{pearl} - q$ string as a $q - g - \bar{q}$ string using existing fragmentation mechanism in PYTHIA

> Fragment the $q - g_{pearl} - \bar{q}$ string system from the \bar{q} end, reversing the hadron IDs > Pick up quark and energy from p_{pearl} for "free" when stepping over junction

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